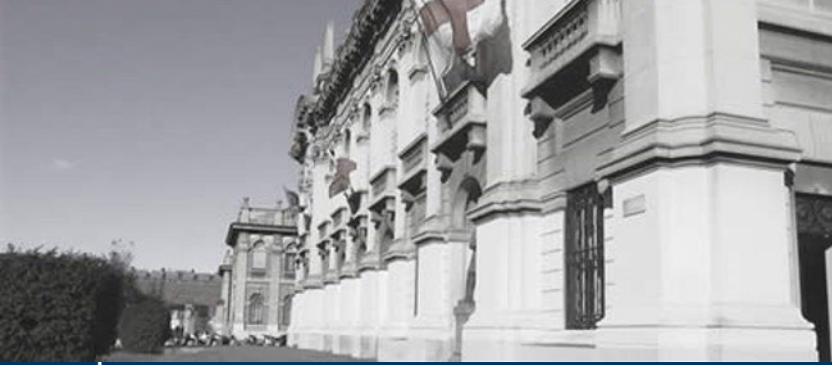




# Electronics – 96032

 POLITECNICO DI MILANO



## Laplace Transforms, Linear Circuits and Bode Plots

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Slides are supplementary material and are NOT a replacement for textbooks and/or lecture notes



- Laplace transforms
- Application to LTI circuits
- Bode plots
- Appendix: Complex poles



$$F(s) = \int_0^{+\infty} f(t)e^{-st}dt = \mathcal{L}(f)$$

$\mathcal{L}$  is a **linear** operator:

$$\mathcal{L}(\alpha f(t) + \beta g(t)) = \alpha \mathcal{L}(f) + \beta \mathcal{L}(g)$$

Note: in the following, all signals are defined only for  $t \geq 0$ , i.e., they are (implicitly) multiplied by the unit step function  $u(t)$



## 1. Time differentiation:

$$\mathcal{L}(f'(t)) = sF(s) - f(0)$$

- If we think in terms of **distributional derivatives** (e.g., Dirac delta as derivative of step function), then  $f(0) \rightarrow f(0^-)$  (= 0 in our case)
- If we think in terms of **classical derivatives** (e.g., zero as derivative of step function), then  $f(0) \rightarrow f(0^+)$



## 2. “Frequency” differentiation:

$$\mathcal{L}(tf(t)) = -\frac{dF(s)}{ds}$$

- More generally

$$\mathcal{L}(t^n f(t)) = (-1)^n \frac{d^n F(s)}{ds^n}$$



3. Time integration:

$$\mathcal{L}\left(\int_0^t f(\tau)d\tau\right) = \frac{F(s)}{s}$$

4. Frequency integration:

$$\mathcal{L}\left(\frac{f(t)}{t}\right) = \int_s^\infty F(\sigma)d\sigma$$



5. Time shifting:

$$\mathcal{L}(f(t - T)) = e^{-sT} F(s)$$

6. Frequency shifting:

$$\mathcal{L}(e^{at} f(t)) = F(s - a)$$

7. Convolution:

$$\mathcal{L}(f(t) * g(t)) = F(s)G(s)$$

- Dual properties is more complicated, involving integral in complex domain



8. Initial value theorem:

$$f(0^+) = \lim_{s \rightarrow \infty} sF(s)$$

9. Final value theorem:

$$f(+\infty) = \lim_{s \rightarrow 0} sF(s)$$

Both can be extended to compute the values of the derivatives of  $f$  just by recalling #1. For first derivatives, this means

$$f'(0^+) = \lim_{s \rightarrow \infty} s^2 F(s) \quad f'(+\infty) = \lim_{s \rightarrow 0} s^2 F(s)$$



## 10. Scaling

$$\mathcal{L}(f(at)) = \frac{1}{|a|} F\left(\frac{s}{a}\right)$$

In particular, for  $a = -1$  (time reversal) we get:

$$\mathcal{L}(f(-t)) = F(-s)$$



- Dirac delta function:

$$\mathcal{L}(\delta(t)) = \int_0^{+\infty} \delta(t) e^{-st} dt = 1$$

- Heaviside step function (integral of  $\delta(t)$ ):

$$\mathcal{L}(u(t)) = \frac{1}{s}$$

- Ramp function (integral of  $u(t)$ )

$$\mathcal{L}(t) = \frac{1}{s^2} \quad \text{Also obtained from \#2}$$



- Rectangle function:

$$\mathcal{L}(\text{rect}(0, T)) = \mathcal{L}(u(t) - u(t - T)) = \frac{1 - e^{-sT}}{s}$$

- Decreasing exponential function (from #6, with  $f(t) = u(t)$  and  $a = -1/\tau$ ):

$$\mathcal{L}(e^{-t/\tau}u(t)) = \frac{1}{s + 1/\tau} = \frac{\tau}{1 + s\tau}$$



# Differentiation example

13

$$f(t) = e^{-t/\tau} \Rightarrow F(s) = \frac{\tau}{1 + s\tau} \quad \mathcal{L}(f'(t)) = sF(s) = \frac{s\tau}{1 + s\tau}$$

In fact, in a distribution frame:

$$f'(t) = \delta(t) - \frac{1}{\tau} e^{-t/\tau}$$

$$\mathcal{L}(f'(t)) = 1 - \frac{1}{\tau} \frac{\tau}{1 + s\tau} = \frac{s\tau}{1 + s\tau}$$



- Sine function

$$\mathcal{L}(\sin(\omega t)) = \frac{\omega}{s^2 + \omega^2}$$

- Cosine function (from #1)

$$\mathcal{L}(\cos(\omega t)) = \frac{s}{s^2 + \omega^2}$$

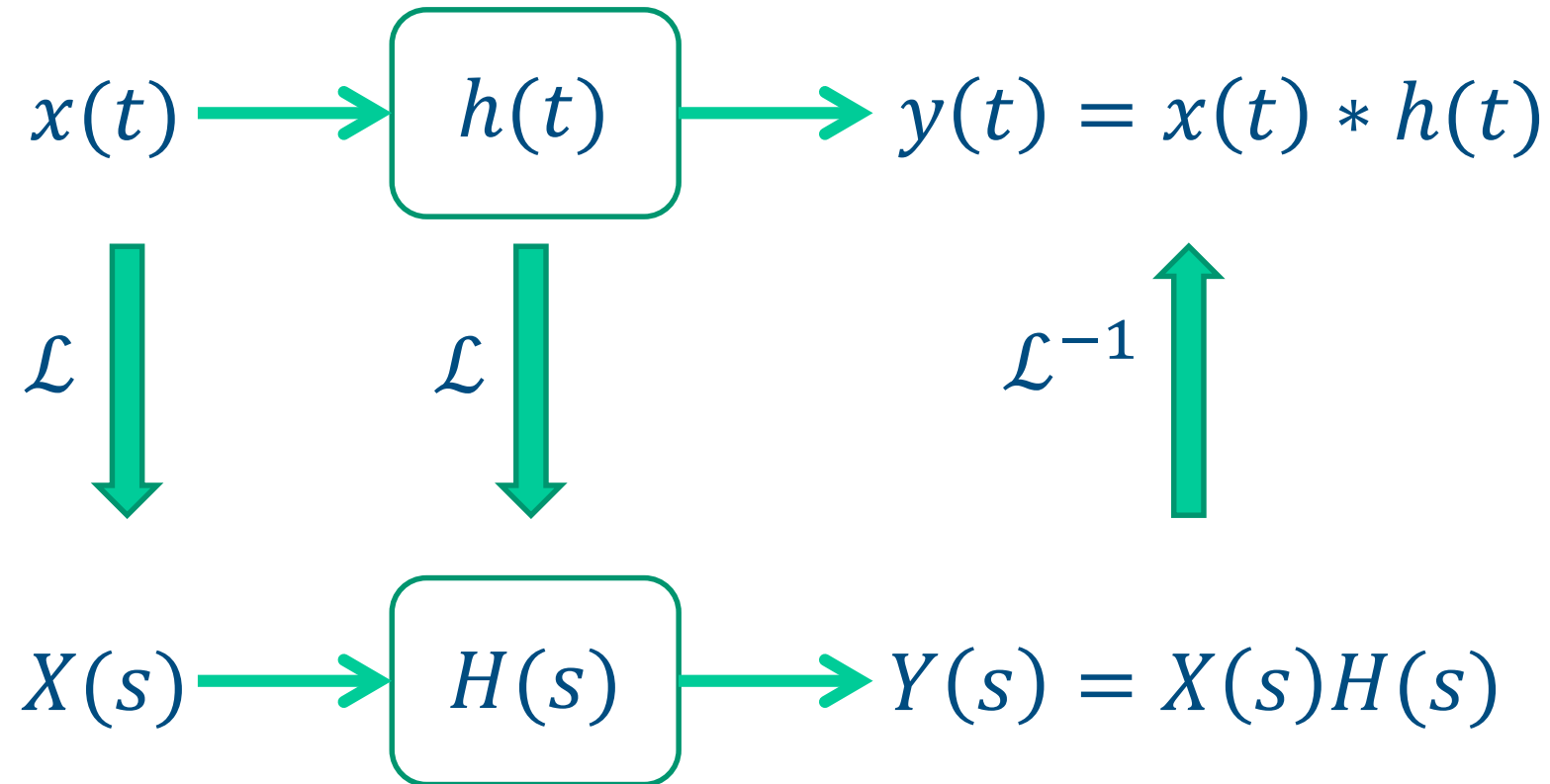


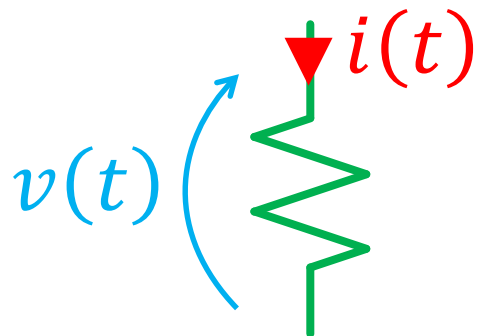
- Laplace transforms
- Application to LTI circuits
- Bode plots
- Appendix: Complex poles



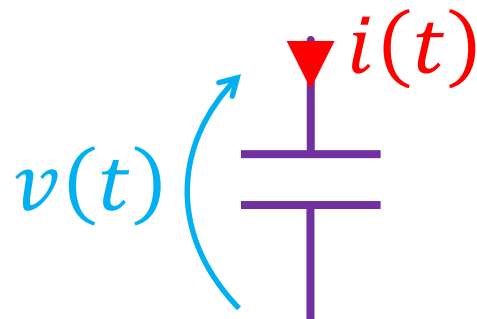
# Application to LTI systems

16

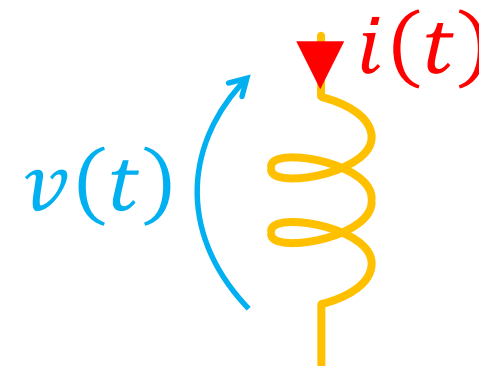




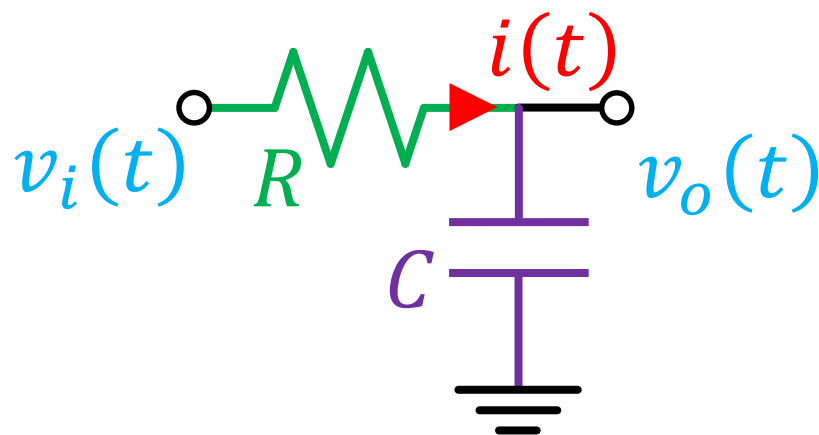
$$v(t) = Ri(t)$$
$$V(s) = RI(s)$$
$$\frac{V(s)}{I(s)} = R$$



$$i(t) = C \frac{dv(t)}{dt}$$
$$I(s) = sCV(s)$$
$$\frac{V(s)}{I(s)} = \frac{1}{sC} = Z_C$$



$$v(t) = L \frac{di(t)}{dt}$$
$$V(s) = sLI(s)$$
$$\frac{V(s)}{I(s)} = sL = Z_L$$



$$\begin{cases} v_i = v_o + Ri \\ i = C \frac{dv_o}{dt} \end{cases} \Rightarrow RC \frac{dv_o}{dt} + v_o = v_i$$

$$RCz + 1 = 0 \Rightarrow z = -\frac{1}{RC} \Rightarrow v_o = K e^{-\frac{t}{RC}} \quad (\text{homogeneous sol.})$$

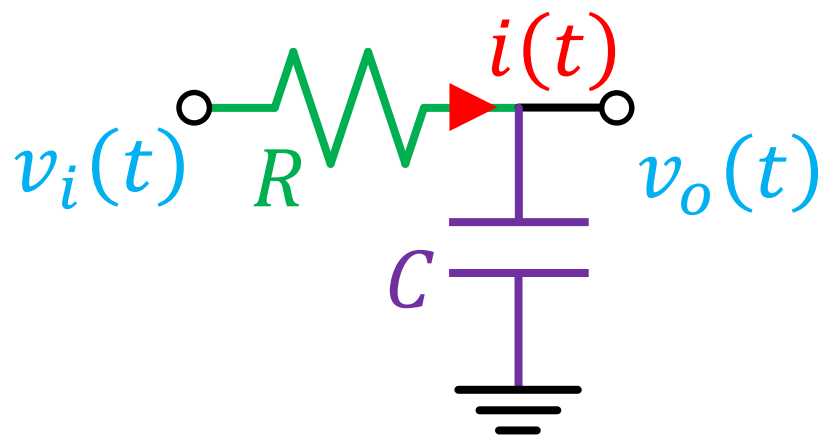
$$v_i = Au(t) \Rightarrow v_o = A \quad (\text{particular solution})$$

$$v_o(0) = 0 \Rightarrow K = -A \Rightarrow v_o = A \left( 1 - e^{-\frac{t}{RC}} \right)$$



# RC network – Laplace domain

19



Single, real pole  
in  $s = -1/CR$

$$\frac{V_o(s)}{V_i(s)} = \frac{1/sC}{R + 1/sC} = \frac{1}{1 + sCR}$$

Poles are the eigenvalues of the characteristic equation  $\Rightarrow$  they represent the time constants of the output



- Delta function response,  $v_i(t) = A\delta(t)$

$$V_i(s) = A \Rightarrow V_o = \frac{A}{1 + sCR} = \frac{A}{\tau} \frac{\tau}{1 + s\tau} \Rightarrow v_o = \frac{A}{\tau} e^{-t/\tau}$$

- Step function response,  $v_i(t) = Au(t)$

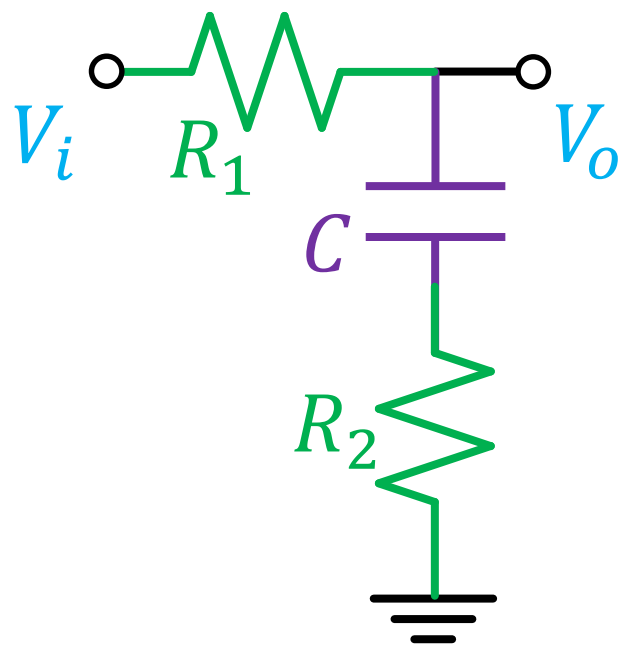
$$V_i(s) = \frac{A}{s} \Rightarrow V_o = \frac{A}{1 + s\tau} \frac{1}{s} = A \left( \frac{1}{s} - \frac{\tau}{1 + s\tau} \right) \Rightarrow$$

$$v_o = A \left( u(t) - e^{-t/\tau} u(t) \right) = A(1 - e^{-t/\tau})u(t)$$



## Ex: Lag network (step response)

21



$$\frac{V_o}{V_i} = \frac{R_2 + 1/sC}{R_1 + R_2 + 1/sC} = \frac{1 + sCR_2}{1 + sC(R_1 + R_2)}$$

$$V_o = \frac{1 + sCR_2}{1 + sC(R_1 + R_2)} \frac{A}{s}$$

$$V_o = A \left( \frac{1}{s} - \frac{CR_1}{1 + sC(R_1 + R_2)} \right) \Rightarrow v_o(t) = A \left( 1 - \frac{R_1}{R_1 + R_2} e^{-t/\tau} \right)$$



# Alternative (rough) method

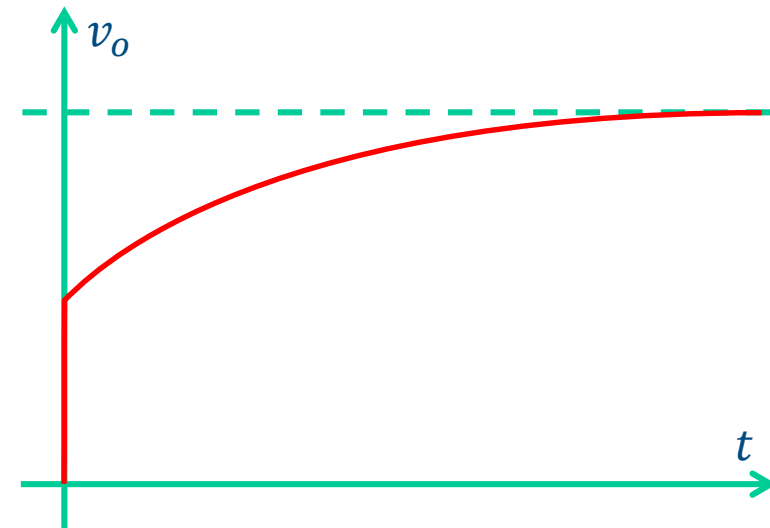
22

1) Compute:

$$v_o(0) = \lim_{s \rightarrow \infty} sV_o(s) = A \frac{R_2}{R_1 + R_2}$$

$$v_o(\infty) = \lim_{s \rightarrow 0} sV_o(s) = A$$

2) Connect extremes with an exponential curve having time constant determined by the single pole





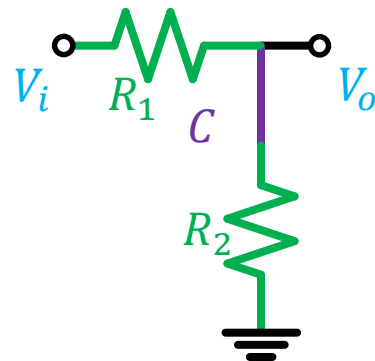
# One more approach

23

Voltage across a capacitor cannot change abruptly (unless you are applying a  $\delta$ -function current)

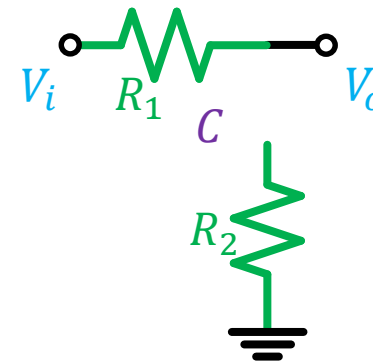
- $v_c(0^+) = v_c(0^-) = 0$

- $v_o(0^+) = \frac{AR_1}{R_1 + R_2}$



For  $t \rightarrow \infty$ , the capacitor behaves as an open circuit

- $v_o(\infty) = A$

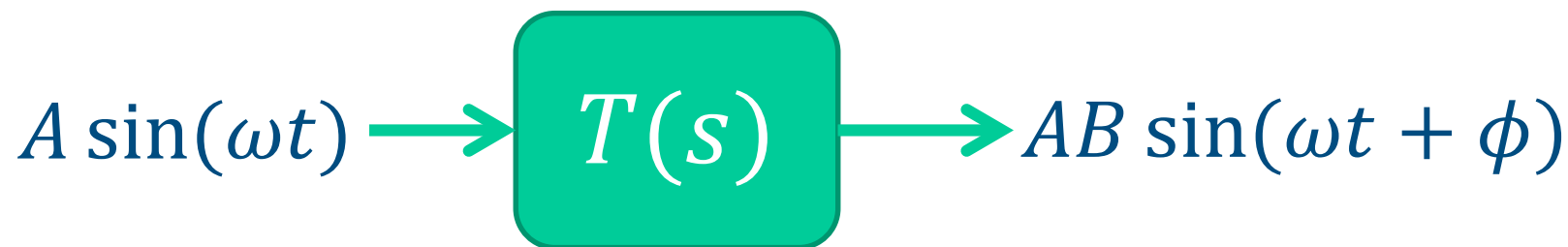




- Laplace transforms
- Application to LTI circuits
- **Bode plots**
- Appendix: Complex poles



- The output of a stable LTI system to a sinusoidal input contains a transient and a steady-state term
- The steady-state term is also sinusoidal at the same frequency as the input, but with different amplitude and phase:



where

$$B = |T(j\omega)| \quad \phi = \angle(T(j\omega))$$



- Provide an efficient way to plot the frequency response (magnitude and phase) of LTI systems
- We consider **asymptotic** Bode plots, i.e., piecewise linear approximations on a suitable scale:
  - Magnitude:  $\text{dB} = 20 \log_{10} |\cdot|$
  - Phase: linear scale
  - Frequency: log scale



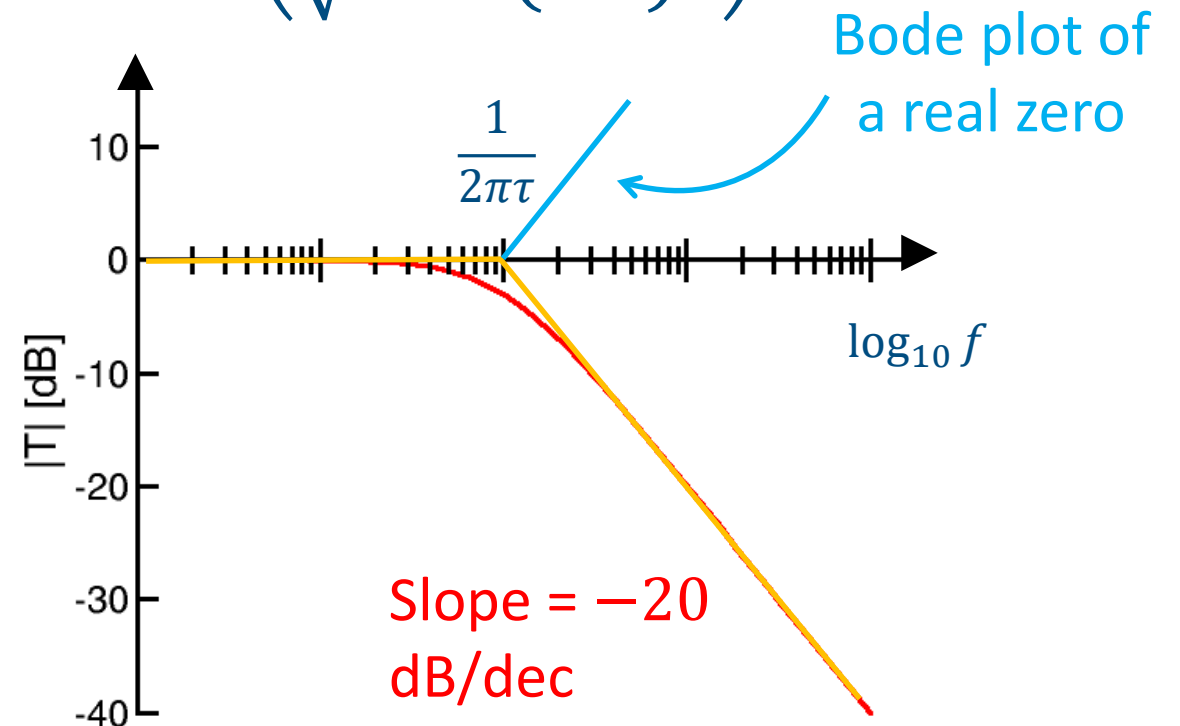
# Single real pole: magnitude

27

$$T(s) = \frac{1}{1 + s\tau} \Rightarrow |T(j\omega)|_{dB} = 20 \log_{10} \left( \frac{1}{\sqrt{1 + (\omega\tau)^2}} \right)$$

$$= -20 \log_{10} \sqrt{1 + (\omega\tau)^2}$$

$$\approx \begin{cases} 0 & \omega\tau \ll 1 \\ -20 \log_{10} \omega\tau & \omega\tau \gg 1 \end{cases}$$

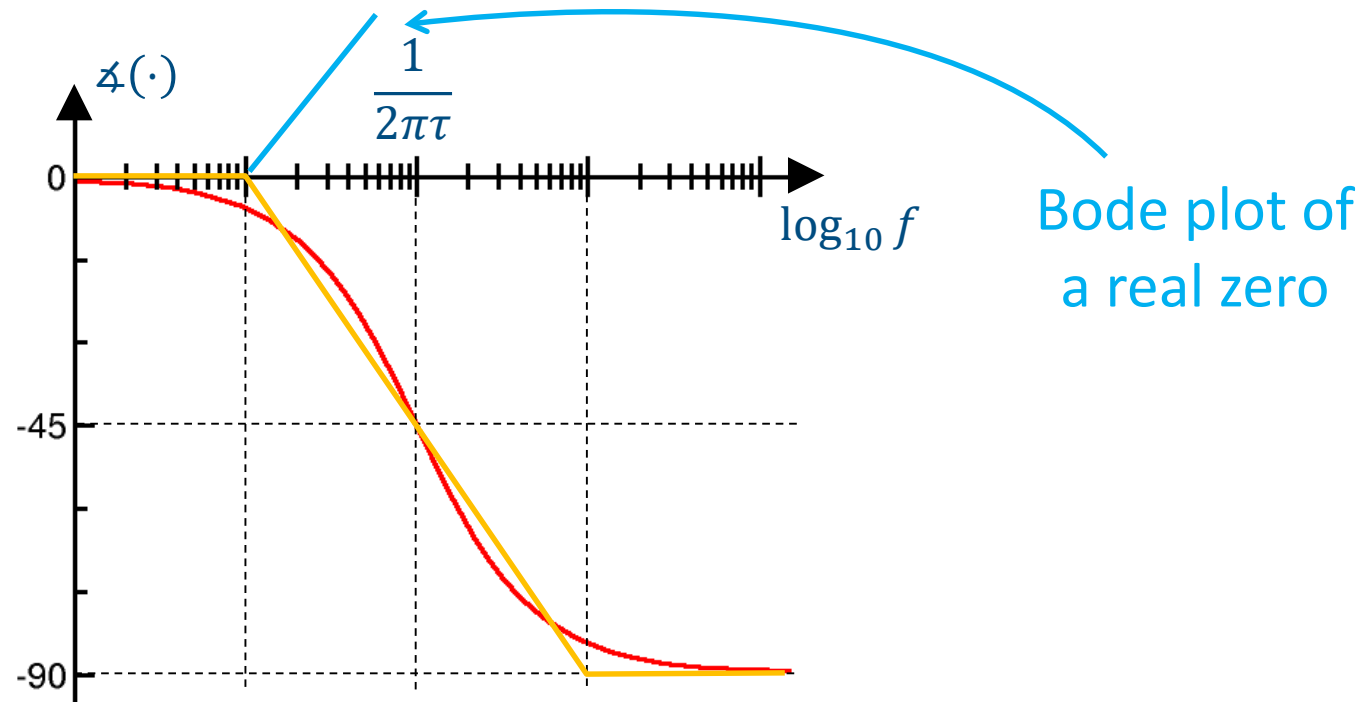




# Single real pole: phase

28

$$\angle(T(j\omega)) = \angle\left(\frac{1}{1 + j\omega\tau}\right) = -\arctan(\omega\tau)$$





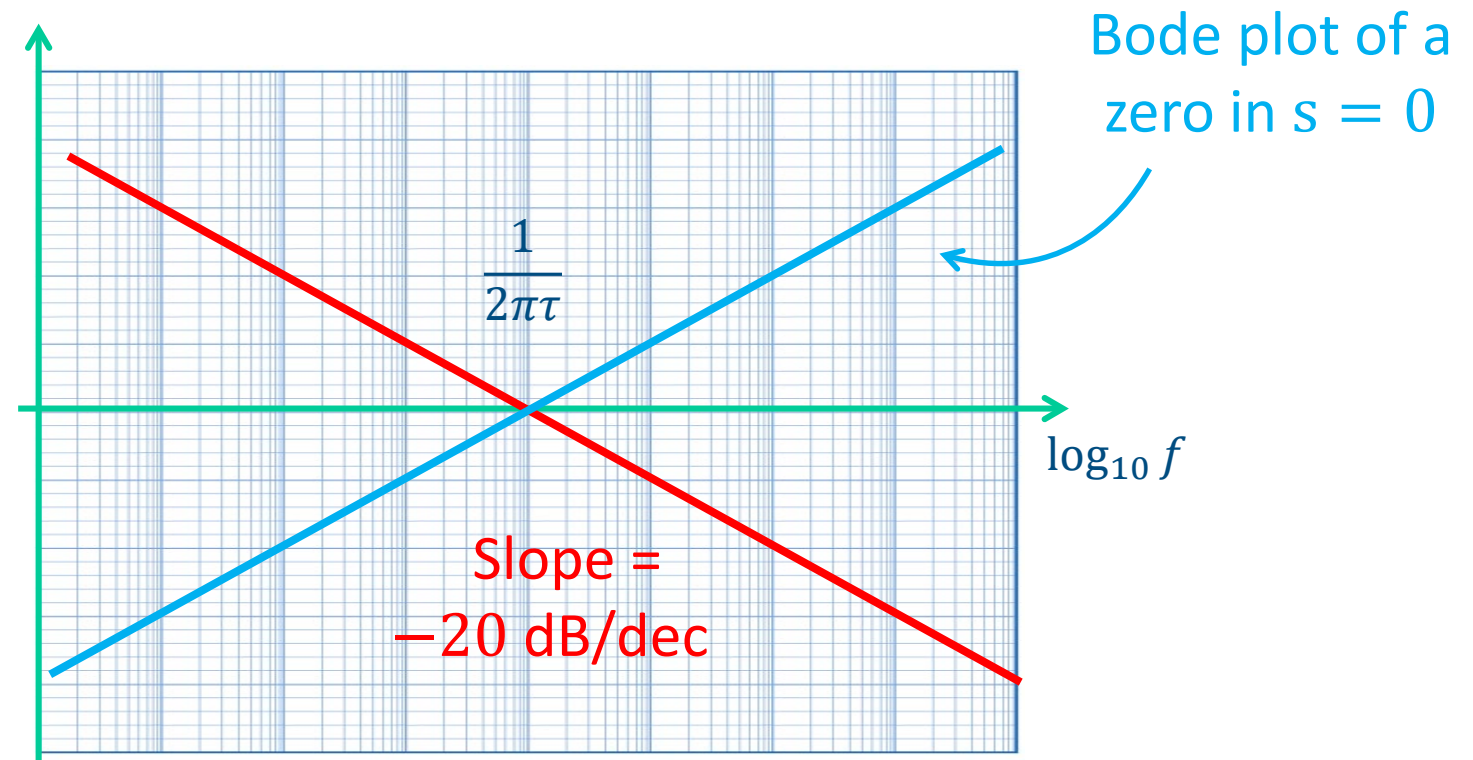
# Pole (zero) in $s = 0$

29

$$T(s) = \frac{1}{s\tau} \Rightarrow |T(j\omega)|_{dB} = 20 \log_{10} \left( \frac{1}{\omega\tau} \right) = -20 \log_{10} \omega\tau$$

x-axis intercept:

$$|T(j\omega)| = 1 \Rightarrow \omega = \frac{1}{\tau}$$





# General case (real singularities)

30

- $T(s)$  can always be expressed as

$$T(s) = G \frac{(1 + s\tau_{z1})(1 + s\tau_{z2}) \dots (1 + s\tau_{zn})}{(1 + s\tau_{p1})(1 + s\tau_{p2}) \dots (1 + s\tau_{pm})}$$

- Thanks to log and arctan properties:

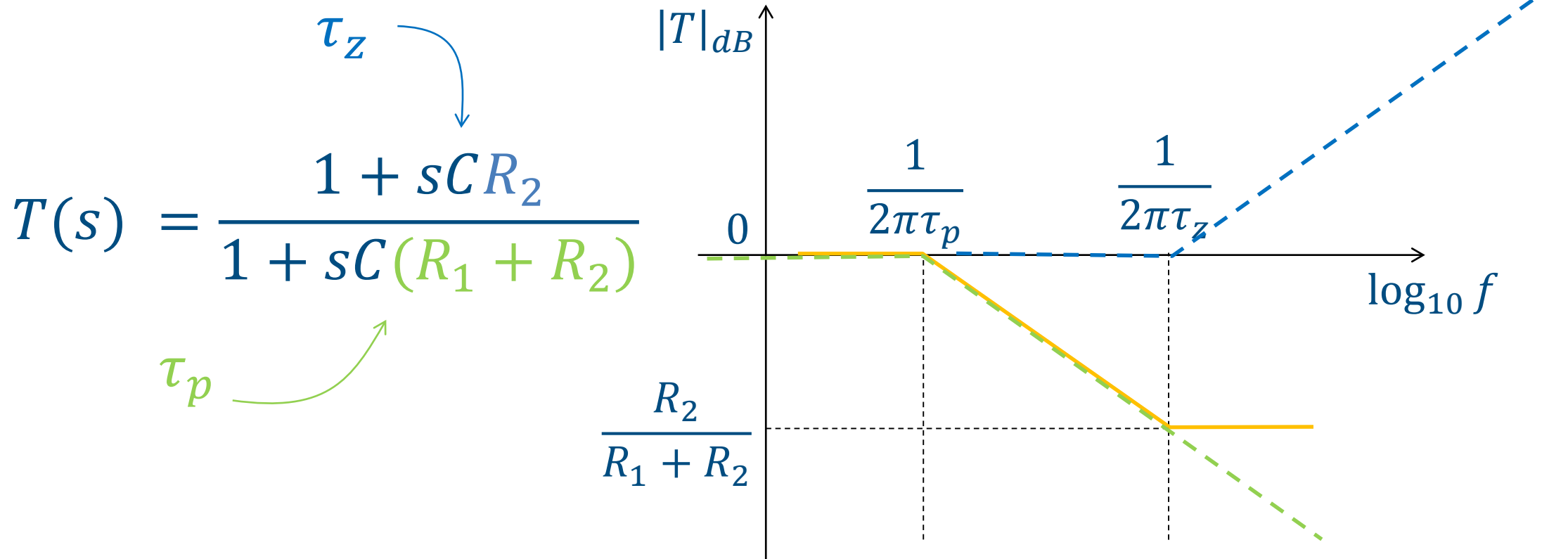
$$|T|_{dB} = G_{dB} + \sum_{i=1}^n |1 + j\omega\tau_{zi}|_{dB} - \sum_{j=1}^m |1 + j\omega\tau_{pj}|_{dB}$$

$$\angle(T) = \sum_{i=1}^n \angle(1 + j\omega\tau_{zi}) - \sum_{j=1}^m \angle(1 + j\omega\tau_{pj})$$



# Ex: Lag network

31



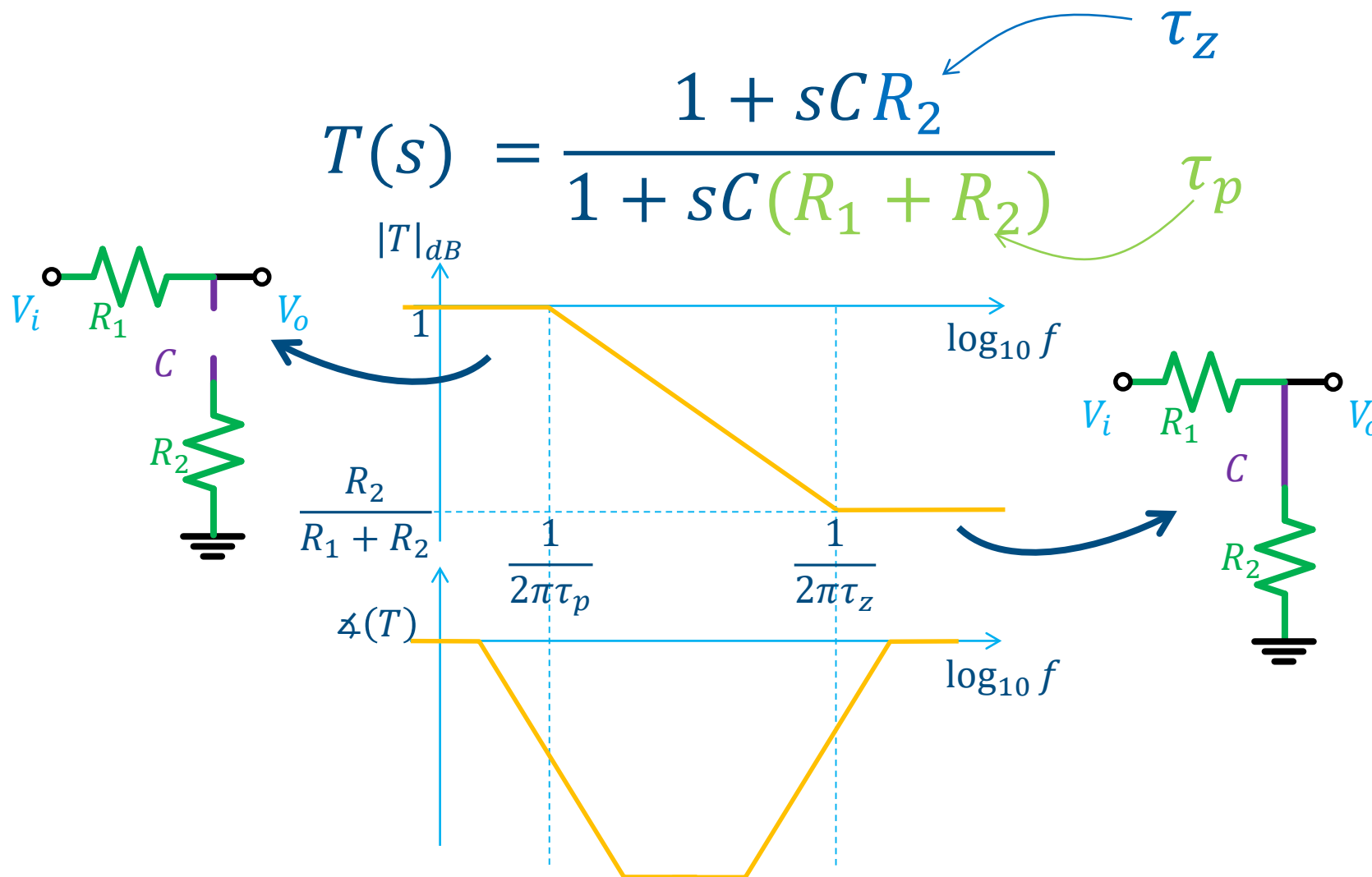


- Magnitude:
  - At every **zero** frequency, the slope is **increased** by 1 (20 dB/dec)
  - At every **pole** frequency, the slope is **reduced** by 1 (20 dB/dec)
- Phase:
  - Composition is a bit more complicated, but a rough view can be obtained by abruptly adding  $\pm 90^\circ$  at every zero/pole frequency



# Asymptotic values

33





- Laplace transforms
- Application to LTI circuits
- Bode plots
- **Appendix: Complex poles**



- We consider a second order LPF having

$$T(s) = \frac{\omega_p^2}{s^2 + 2\xi\omega_p s + \omega_p^2} = \frac{\omega_p^2}{s^2 + s\left(\frac{\omega_p}{Q}\right) + \omega_p^2} \Rightarrow Q = \frac{1}{2\xi}$$

where

$\xi$  = damping factor (mostly used in system control)

$Q$  = quality factor (mostly used in electronics/filter design)



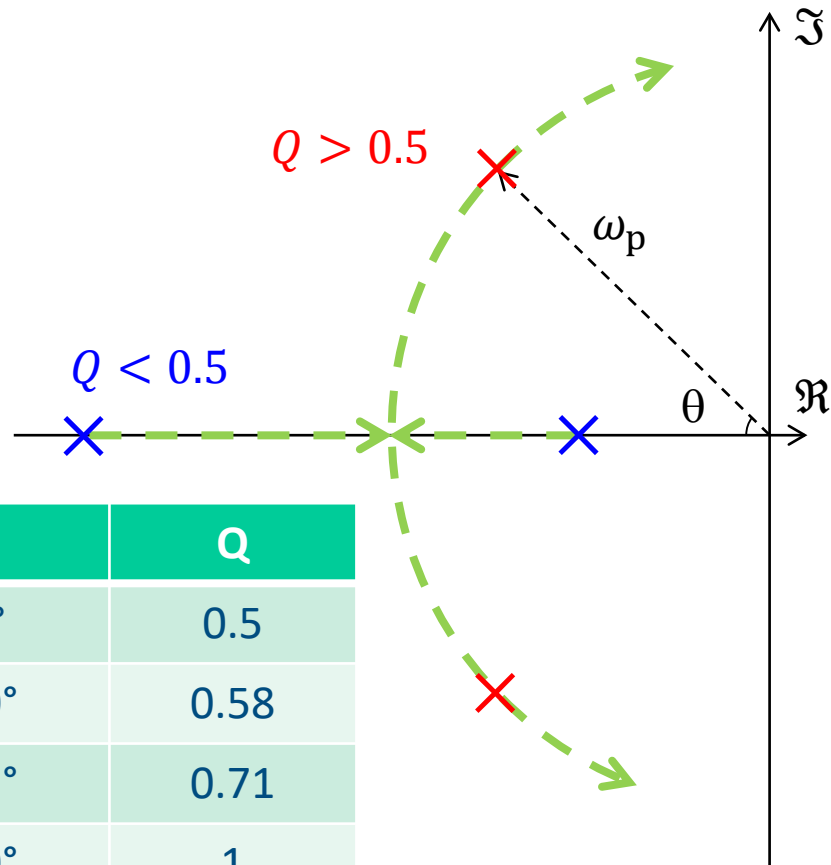
- The roots are

$$\omega = -\frac{\omega_p}{2Q} \pm \omega_p \sqrt{\frac{1}{4Q^2} - 1} = -\frac{\omega_p}{2Q} \left(1 \pm \sqrt{1 - 4Q^2}\right)$$

$Q < 0.5 \Rightarrow$  poles are real and separated

$Q = 0.5 \Rightarrow$  poles are real and coincident

$Q > 0.5 \Rightarrow$  poles are complex conjugated



| $\theta$ | $Q$  |
|----------|------|
| 0°       | 0.5  |
| 30°      | 0.58 |
| 45°      | 0.71 |
| 60°      | 1    |
| 75°      | 1.93 |

• For  $Q > 0.5$  we have

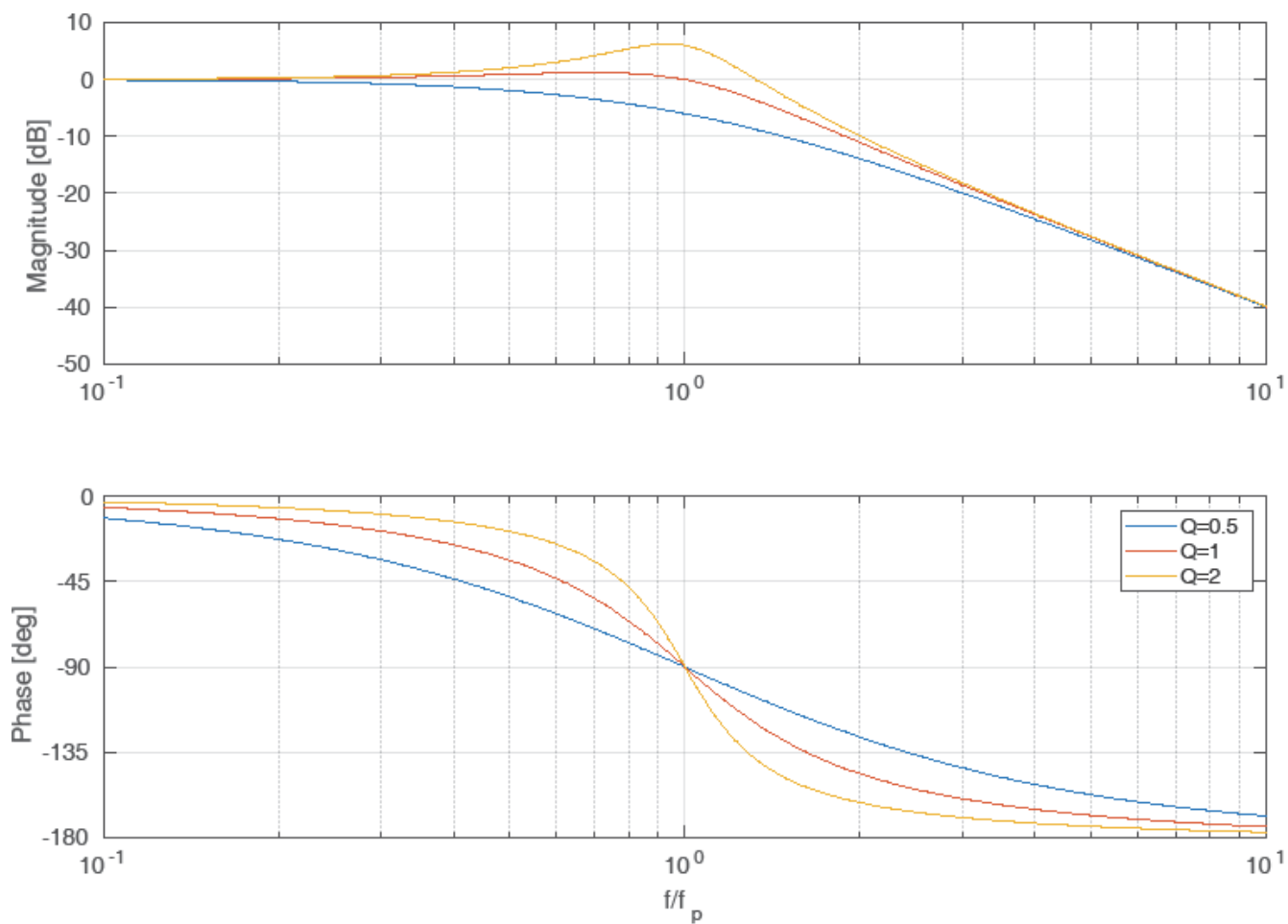
- A real part equal to  $-\frac{\omega_p}{2Q}$
- The frequency of the poles is  $\omega_p/2\pi$
- At  $\omega = \omega_p$  we have

$$T(j\omega_p) = \frac{\omega_p^2}{(j\omega_p)^2 + j\omega_p \left(\frac{\omega_p}{Q}\right) + \omega_p^2} = -jQ$$
$$|T(j\omega_p)| = Q$$



# Bode diagrams ( $Q \geq 0.5$ )

38





- The TF of a second-order BPF is

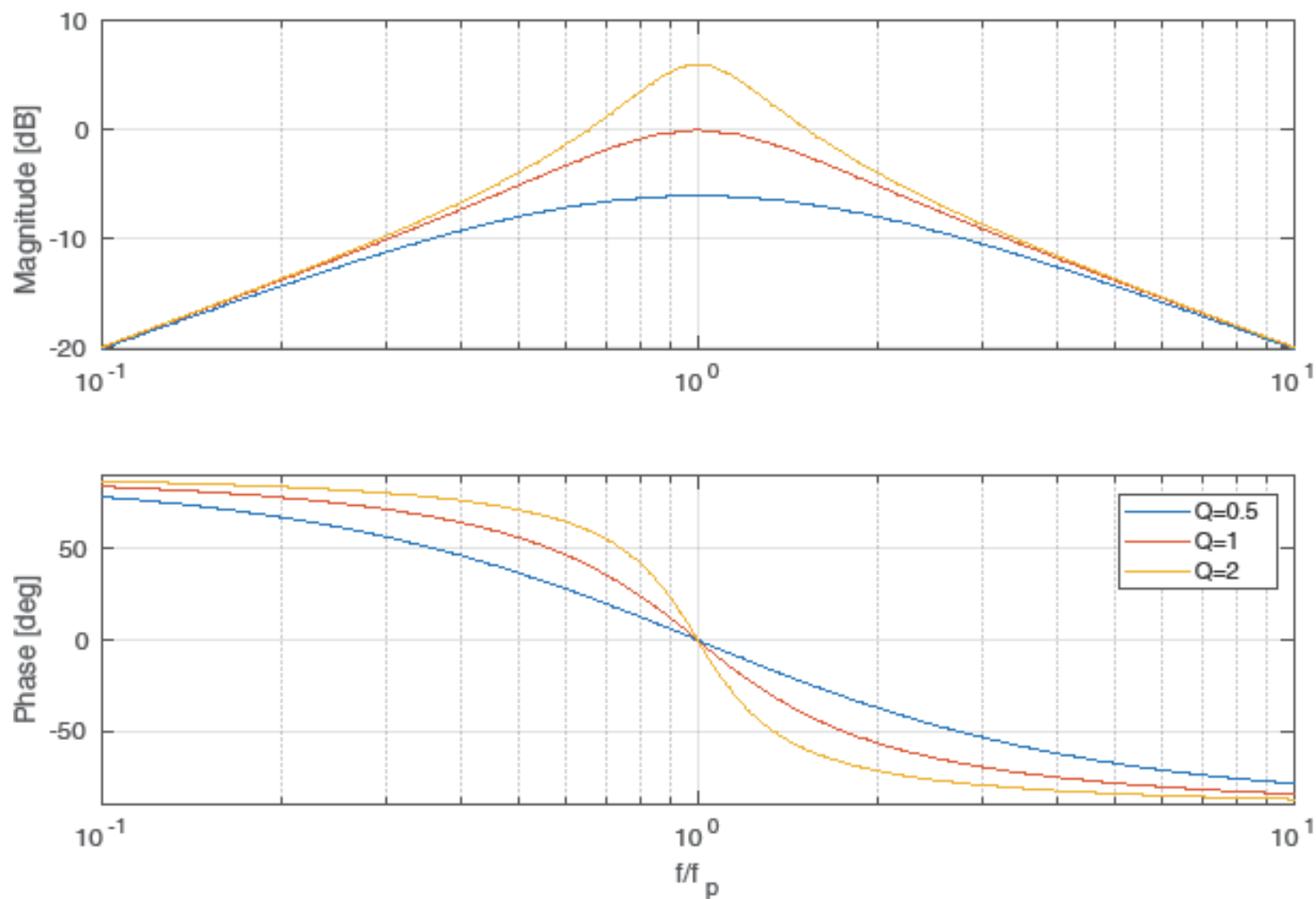
$$T(s) = \frac{s\omega_p}{s^2 + s\left(\frac{\omega_p}{Q}\right) + \omega_p^2}$$

- Pole position is the same as the LPF's
- We still have  $|T(j\omega_p)| = Q$
- The -3dB BW is  $\frac{f_p}{Q}$



# Bode diagrams ( $Q \geq 0.5$ )

40





1. Compute the step response of a CR filter
2. Design a simple network (made with resistors and capacitors only) having a zero and a pole with  $f_z < f_p$  (a lead network)
3. Consider a BP filter with real poles, having

$$T(s) = \frac{As\tau_1}{(1 + s\tau_1)(1 + s\tau_2)} \quad (\tau_1 \gg \tau_2)$$

Draw its quantitative Bode diagram and work out the step response